The Johns Hopkins Turbulence Databases (*JHTDB*)

FORCED ISOTROPIC TURBULENCE DATA SET (EXTENDED)

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The data is from a direct numerical simulation of forced isotropic turbulence on a 1024^3 periodic grid, using a pseudo-spectral parallel code. Time integration of the viscous term is done analytically using integrating factor. The other terms are integrated using a second-order Adams-Bashforth scheme and the nonlinear term is written in vorticity form[1]. The simulation is de-aliased using phase-shift and a $2\sqrt{2}/3$ truncation[2,3]. Energy is injected by keeping constant the total energy in modes such that their wave-number magnitude is less or equal to 2. After the simulation has reached a statistical stationary state, 5028 frames of data, which includes the 3 components of the velocity vector and the pressure, are generated and ingested into the database. The duration of the stored data is about five large-eddy turnover times.

Simulation parameters:
Domain: 2π x 2π x 2π (i.e. range of x_1 , x_2 and x_3 is [0,2π])
Grid: 1024^3
Viscosity (ν) = 0.000185
Simulation time-step Δt = 0.0002
Data are stored separated by δt = 0.002 (i.e. every 10 DNS time-steps is stored)
Time stored: between t=0 and 10.056 (5028 time samples separated by δt)

Statistical characteristics of turbulence, time averaged over t=0 and 10.056:

Total kinetic energy, $E_{tot}=\frac{1}{2}\langle u_i u_i\rangle$: $E_{tot}=0.705$

Dissipation, $\varepsilon=2\nu\langle S_{ij}S_{ij}\rangle$: $\varepsilon=0.103$

Rms velocity, $u'=\sqrt{(2/3)E_{tot}}$: $u'=0.686$

Taylor Micro. Scale $\lambda=\sqrt{15\nu u'^2/\varepsilon}$: $\lambda=0.113$

Taylor-scale Reynolds #, $R_\lambda=u'\lambda/\nu$: $R_\lambda=418$

Kolmogorov time scale $\tau_\eta=\sqrt{\frac{\nu}{\varepsilon}}$: $\tau_\eta=0.0424$

Kolmogorov length scale $\eta=\nu^{3/4}\varepsilon^{-1/4}$: $\eta=0.00280$

Integral scale: $L=\frac{\pi}{2u'^2}\int\frac{E(k)}{k}dk$: $L=1.364$

Large eddy turnover time: $T_L=L/u'$: $T_L=1.99$

Statistical characteristics of turbulence, time averaged over t=0 and 2.048:

Total kinetic energy, $E_{tot}=\frac{1}{2}\langle u_i u_i\rangle$: $\qquad$ $E_{tot}=0.695$

Dissipation, $\varepsilon=2\nu\langle S_{ij}S_{ij}\rangle$: $\qquad$ $\varepsilon=0.0928$

Rms velocity, $u'=\sqrt{(2/3)E_{tot}}$: $\qquad$ $u'=0.681$

Taylor Micro. Scale $\lambda=\sqrt{15\nu u'^2/\varepsilon}$: $\qquad$ $\lambda=0.118$

Taylor-scale Reynolds #, $R_\lambda=u'\lambda/\nu$: $\qquad$ $R_\lambda=433$

Kolmogorov time scale $\quad\tau_\eta=\sqrt{\frac{\nu}{\varepsilon}}$: $\qquad$ $\tau_\eta=0.0446$

Kolmogorov length scale $\eta=\nu^{3/4}\varepsilon^{-1/4}$: $\qquad$ $\eta=0.00287$

Integral scale: $L=\frac{\pi}{2u'^2}\int\frac{E(k)}{k}dk$: $\qquad$ $L=1.376$

Large eddy turnover time: $T_L=L/u'$: $\qquad$ $T_L=2.02$

Figures 1 and 2 show radial kinetic energy spectrum compared to the Kolmogorov inertial range spectrum.Figures 3 and 4 show time-series of E_{tot} and Re_λ. The values corresponding to the data in the database are for t>0 and shown in solid lines.
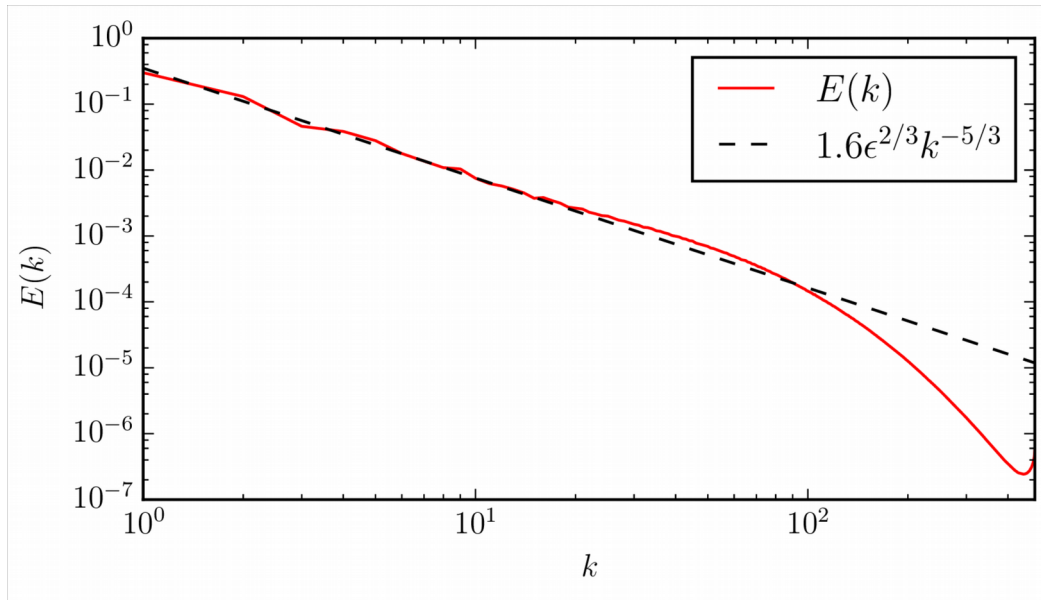

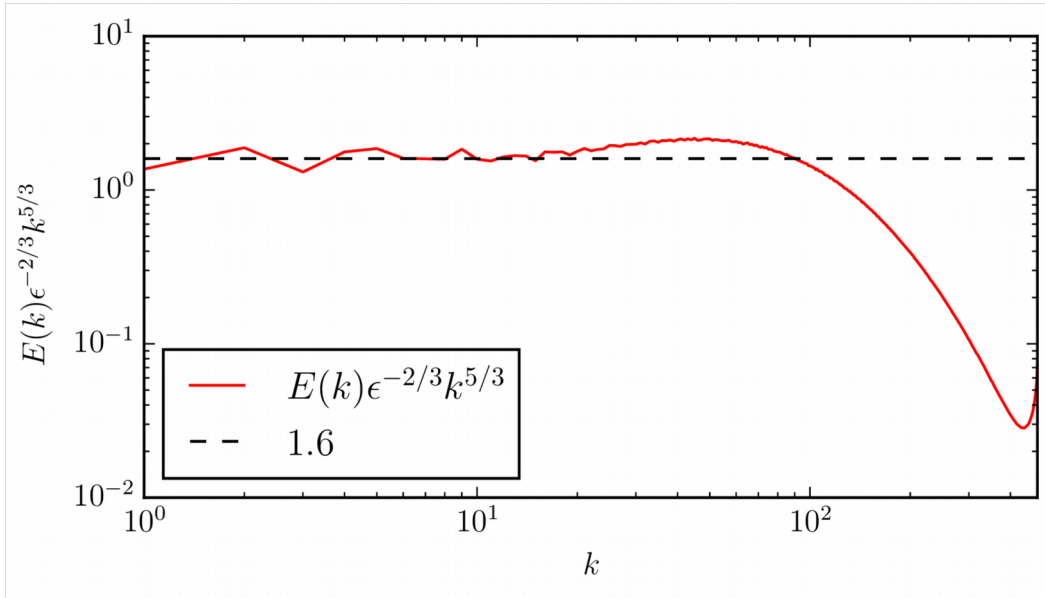Figure 1: Radial kinetic energy spectrum, averaged in time between t=0 and 10.056.

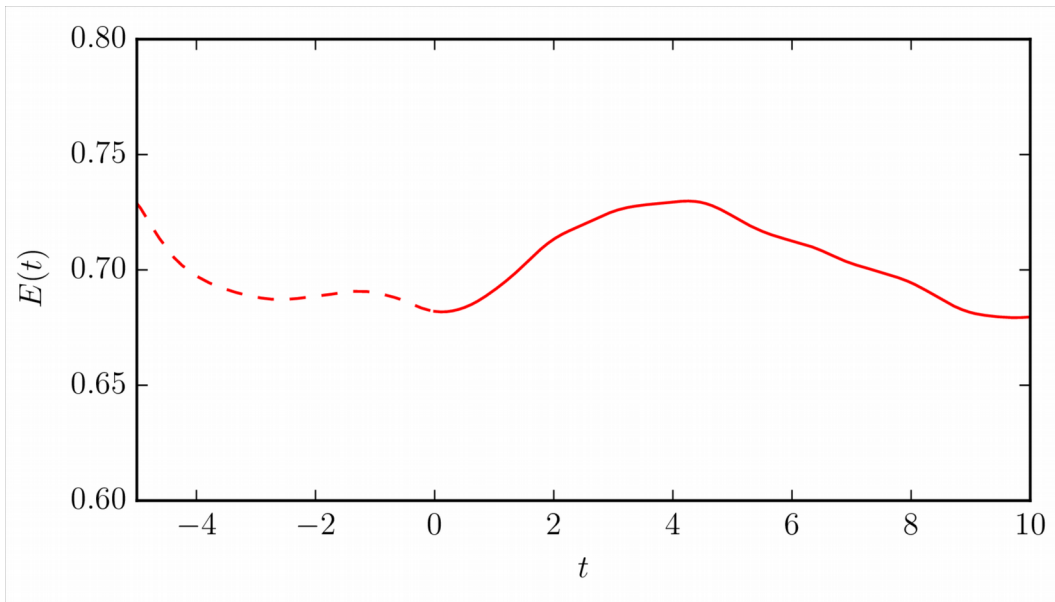Figure 2: Pre-multiplied spectrum, averaged in time between t=0 and 10.056.

Figure 3: Total kinetic energy as function of time. Dashed line is times before ingestion into database. Data corresponding to the database is show using solid line between t=0 and 10.056.
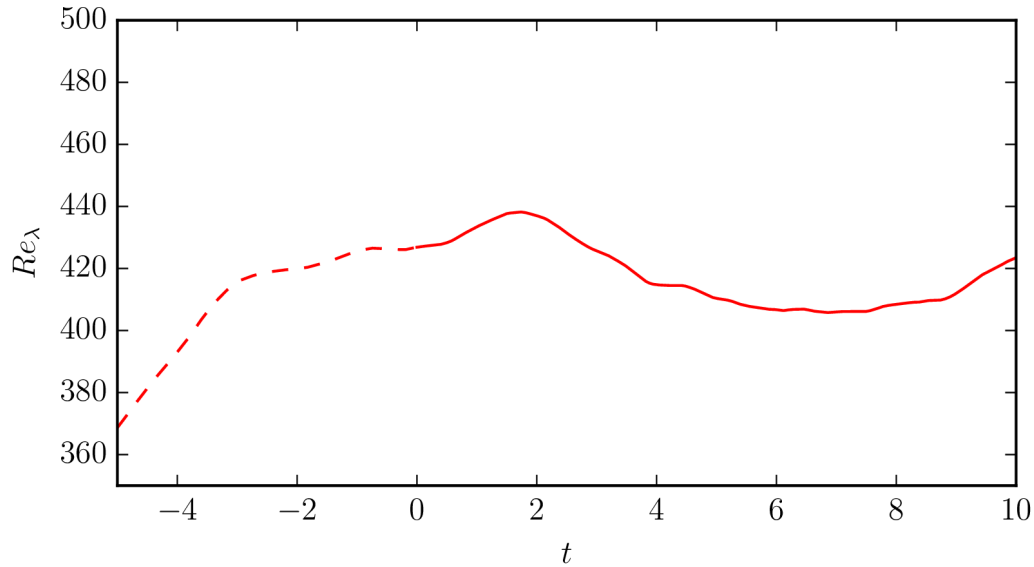
Figure 4: Taylor micro-scale Reynolds number as function of time. Dashed line is times before ingestion into database. Data corresponding to the database is show using solid line between t=0 and 10.056.

References:

1. Cao N.Z. and Chen S.Y., "Statistics and structures of pressure in isotropic turbulence". Phys. Fluids. **11**.2235-2250 (1999)
2. Patterson G.S. and Orszag S.A., "Spectral calculations of isotropic turbulence: efficient removal of aliasing interactions" Phys. Fluids. **14**,2538-2541 (1971).
3. Note: The divergence-free condition in the simulation is enforced based on the spectral representation of the derivatives. The JHTDB analysis tools for gradients are based on finite differencing of various orders. Therefore, when evaluating the divergence using these spatially more localized derivative operators, a non-negligible error in the divergence is obtained, as expected.